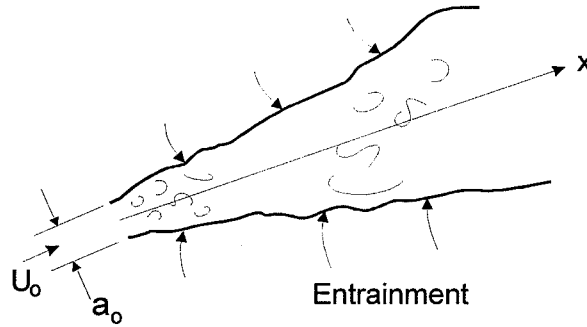


TURBULENT BUOYANT JETS AND PLUMES

Gerhard H. Jirka, University of Karlsruhe

Active dispersal through induced turbulence

Pure jet



Momentum flux

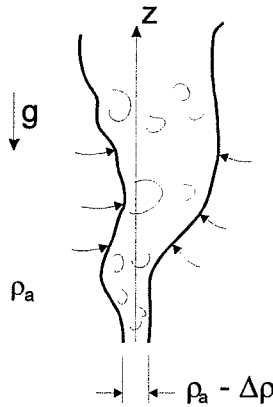
$$M_0 = U_0^2 a_0$$

$$a_0 = D^2 \pi / 4$$

Fully turbulent if

$$Re = \frac{U_0 D}{\nu} \geq 2000$$

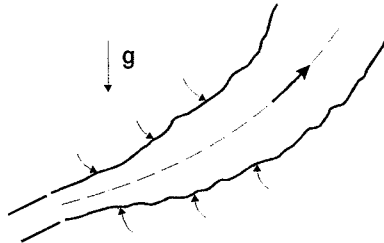
Pure plume



Buoyancy flux

$$J_0 = \left(\frac{\Delta\rho}{\rho_a} g \right) U_0 a_0$$

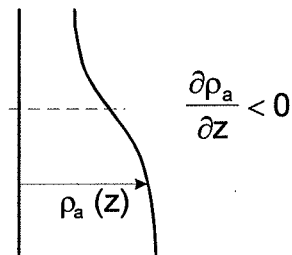
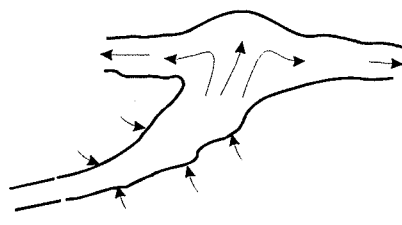
Combination: buoyant jet or forced plume



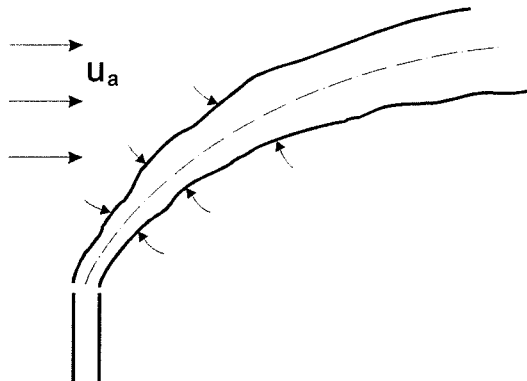
Homogeneous ambient

$$\rho_a = \text{const.}$$

Ambient density stratification



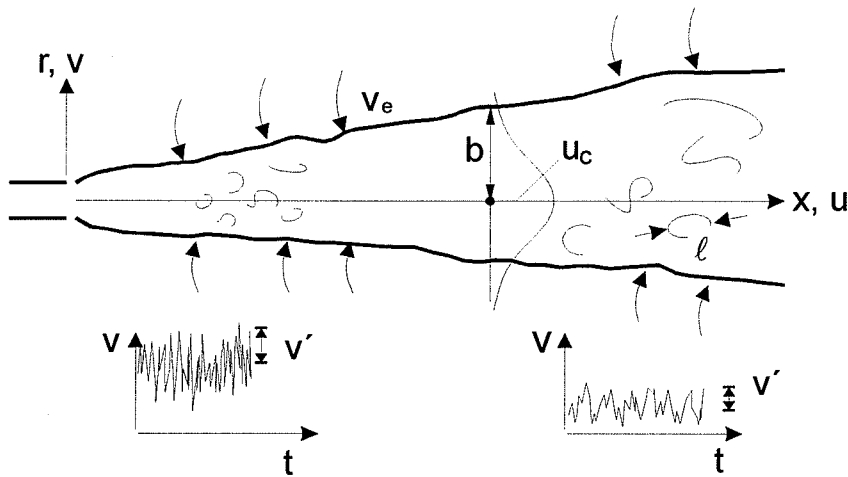
Ambient crossflow



Geometry: 3-D round jet
 multiport diffuser (many 3-D round jets)
 2-D plane jet

Fundamentals: "FREE TURBULENCE" (Prandtl, Tollmien, Taylor ...)

Jet:



Mean:

b = jet width

u_c = centerline velocity

Turbulence:

l = integral length scale

v' = integral velocity scale
(r.m.s.)

1. Invariance $\frac{l}{b} = \text{const.} \quad (\approx 0.3)$

2. Erosion = growth rate $\frac{Db}{Dt} \cong u_c \frac{db}{dx} \sim v'$

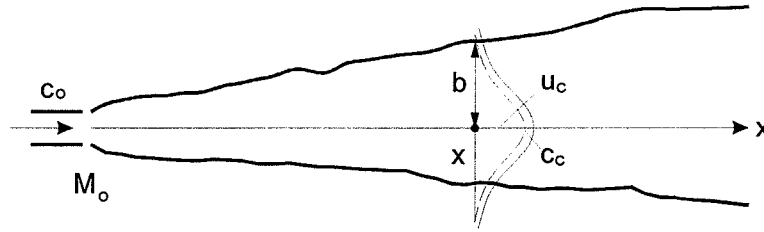
3. Prandtl's approach
(mixing length) $v' = \text{const.} \quad l \frac{\partial u}{\partial r} \sim l \frac{u_c}{b}$

$$u_c \frac{db}{dx} \sim l \frac{u_c}{b} \sim \text{const.} \quad u_c$$

→	$\frac{db}{dx} = k$	<u>linear spreading</u>	<u>"Jet diffusion"</u>	spreading of zone containing momentum (and mass)
	$\frac{v'}{u_c} = \text{const.}$	constant turbulent intensity		
→	$\frac{v_e}{u_c} = \alpha$	<u>constant entrainment rate</u>	<u>"Jet entrainment"</u>	suction of outside fluid into zone (dilution)

Dimensional analysis:

Simple Jet



$$M_0 = U_0^2 a_0 = U_0^2 \frac{\pi D^2}{4}$$

Momentum flux conserved

$$Q_{c_0} = c_0 U_0 a_0$$

Scalar flux

Dynamic jet properties:

$$b = f(x, M_0, a_0, \nu \dots)$$

$$x = [L]$$

$$u_c = g(x, M_0, a_0, \nu \dots)$$

$$M_0 = \left[\frac{L^4}{T^2} \right]$$

$$\frac{b}{x} = \text{const.}, \quad \frac{u_c x}{\sqrt{M_0}} = \text{const.}, \quad \rightarrow \quad b = kx, \quad \frac{u_c}{U_0} = K \frac{\sqrt{a_0}}{x}$$

Experiment: $b = 0.10 x, \quad u_c / U_0 = 7.1 / (x / D)$

Passive properties:

$$Q_{c_0} = c_0 U_0 a_0 \cong c_c u_c b^2 \quad \text{Mass flux conserved}$$

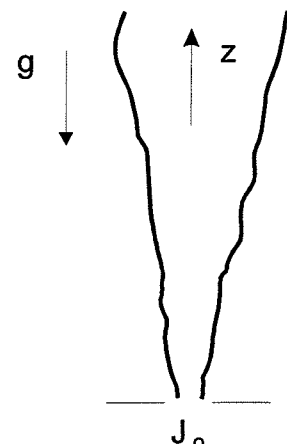
$$\frac{c_c}{c_0} \sim \frac{U_0 a_0}{u_c b^2} = \frac{1}{K} \frac{\sqrt{a_0}}{x}$$

Dilution: $S = \frac{c_0}{c_c} = \bar{K} \frac{x}{\sqrt{a_0}} !$

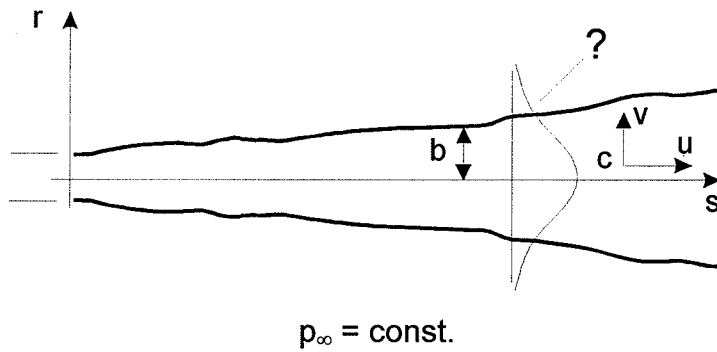
Experiment: $\frac{c_c}{c_0} = \frac{7.1}{x/D}, \quad S_c = 0.14 \frac{x}{D} \quad \text{centerline dilution}$

Simple Plume (J_0, z) $b \sim z, u_c \sim z^{-1/2}, S_c \sim z^{5/3}$

$$J_0 = \left[\frac{L^4}{T^3} \right]$$



Formal solution methods:



Unknowns (u,v,p,c)

Boundary layer approximations: $\frac{b}{s} \ll 1, \frac{\partial}{\partial r} \gg \frac{\partial}{\partial s}$

$$\frac{\partial p}{\partial r} = 0, \frac{\partial p}{\partial s} = 0$$

$$\frac{\partial u}{\partial s} + \frac{1}{r} \frac{\partial rv}{\partial r} = 0 \quad \text{Continuity}$$

$$u \frac{\partial u}{\partial s} + v \frac{\partial u}{\partial r} = -\frac{1}{r} \frac{\partial}{\partial r} (r \overline{u'v'}) \quad \text{Momentum}$$

$$\left[u \frac{\partial c}{\partial s} + v \frac{\partial c}{\partial r} = -\frac{1}{r} \frac{\partial}{\partial r} (r \overline{v'c'}) \right] \quad \text{Scalar transport (uncoupled)}$$

} Jet equation
u,v,c

I.C.: $s = 0: u = U_0, c = C_0, v = 0$

B.C.: $r \rightarrow \infty: u \rightarrow 0, c \rightarrow 0,$
 $\overline{u'v'} \rightarrow 0, \overline{v'c'} \rightarrow 0$

Solutions:

- Similarity methods with simple turbulence closure \Rightarrow classical, simple geometries

$$\overline{u'v'} = \varepsilon \frac{\partial u}{\partial r}$$
- Numerical integration (P.D.E.) with advanced turbulence closure e.g. k- ε \Rightarrow more general geometries

LES
- Integral methods (conversion to O.D.E.)

INTEGRAL METHOD FOR BUOYANT JET ANALYSIS

BUOYANT JET (Round jet, 2-D trajectory, stagnant unstratified ambient)



$$g' = \frac{\Delta\rho}{\rho_a} g$$

1. Profile Specification:

$$\frac{u}{u_c} = e^{-r^2/b^2}$$

$$b = \frac{1}{e} \text{ "Width" (37 \%)}$$

$$\frac{g'}{g'_c} = \frac{c}{c_c} = e^{-\frac{r^2}{(\lambda b)^2}}$$

$$\lambda > 1 \text{ "Dispersion ratio" (turbulent Schmidt number)}$$

2. Definition of integral quantities:

$$\text{Volume flux} \quad Q = 2\pi \int_0^\infty u r \, dr = 2\pi u_c b^2 \underbrace{\int_0^\infty \left(\frac{r}{b}\right) e^{-\left(\frac{r}{b}\right)^2} d\left(\frac{r}{b}\right)}_{1/2} = \pi u_c b^2$$

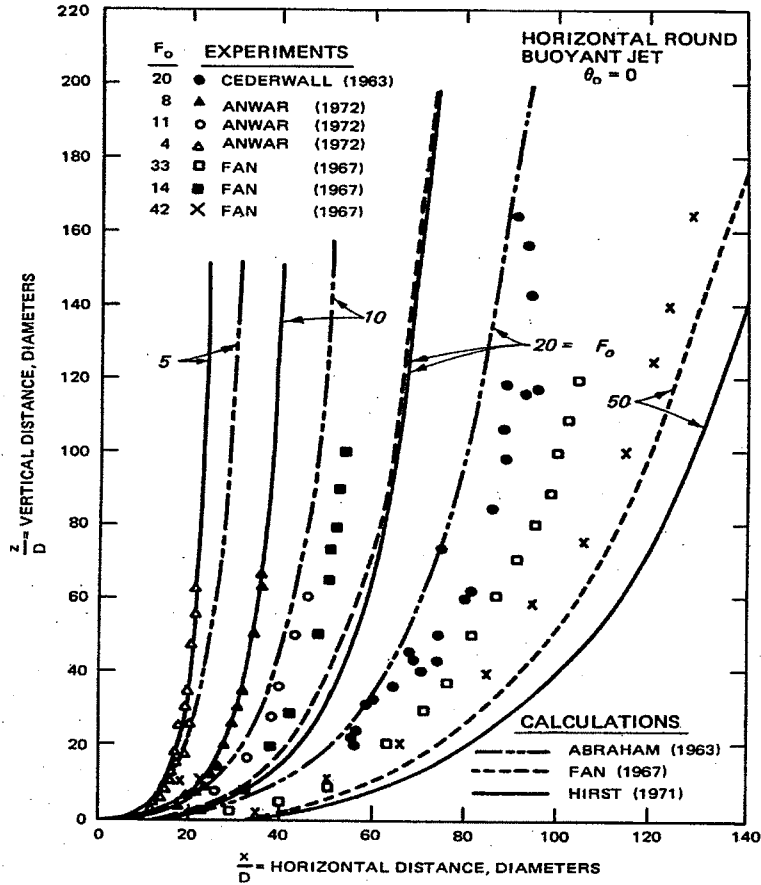
$$\text{Momentum flux} \quad M = \frac{\pi}{2} u_c^2 b^2$$

$$\text{Buoyancy flux} \quad J = \frac{\lambda^2}{1+\lambda^2} \pi u_c g'_c b^2$$

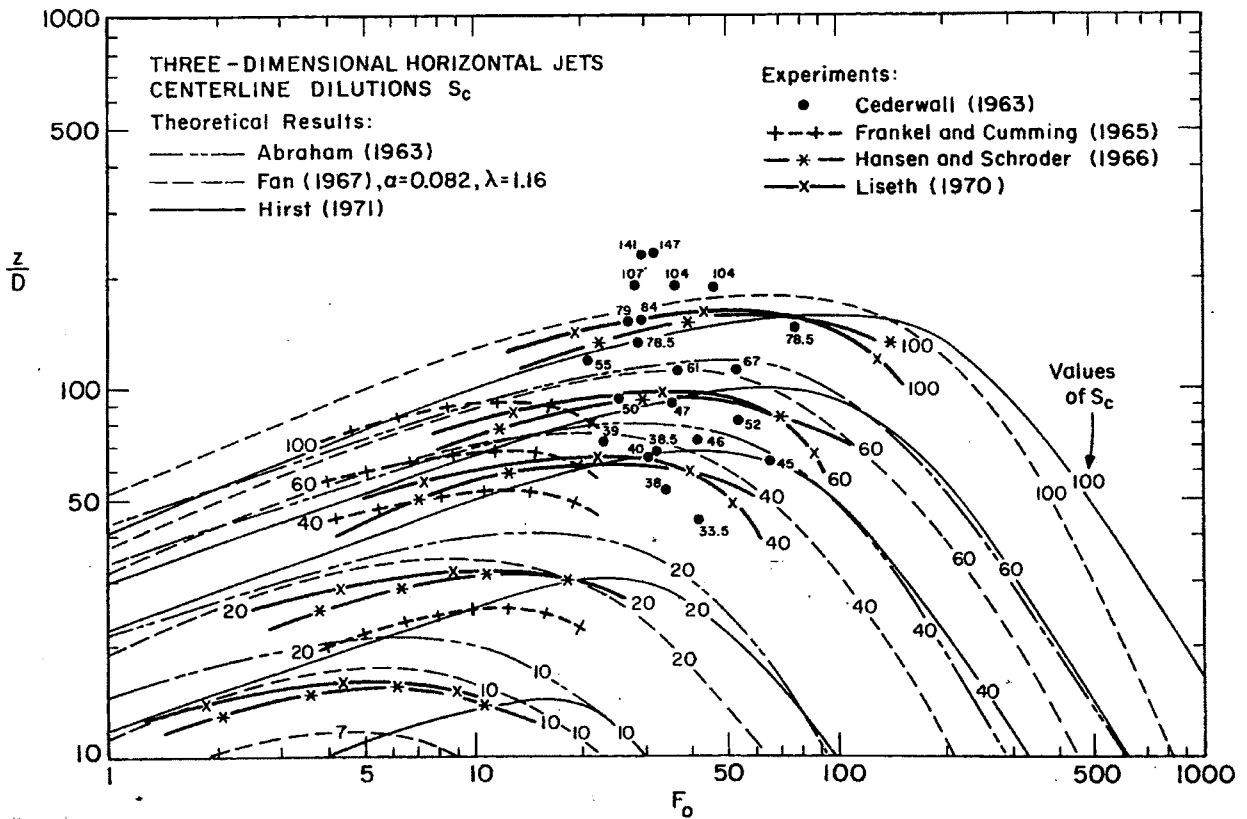
$$\text{Scalar flux} \quad Q_c = \frac{\lambda^2}{1+\lambda^2} \pi u_c c_c b^2$$

(„tracer, pollutant“)

STAGNANT HOMOGENEOUS AMBIENT

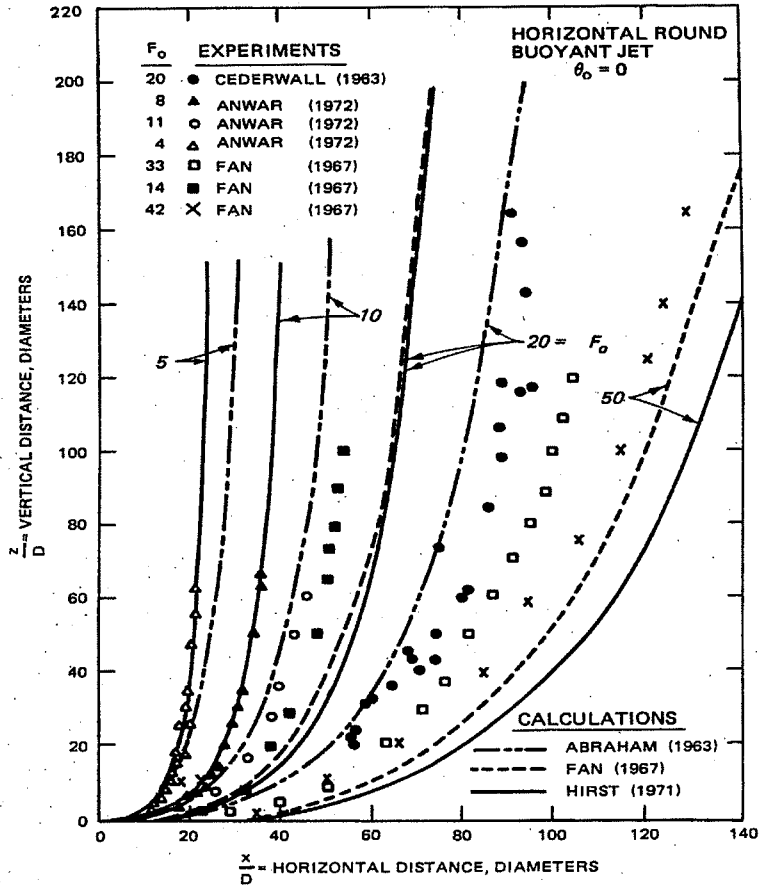


Three-Dimensional Horizontal Jet Trajectories.
Comparison Between Theoretical Predictions and Experimental Data

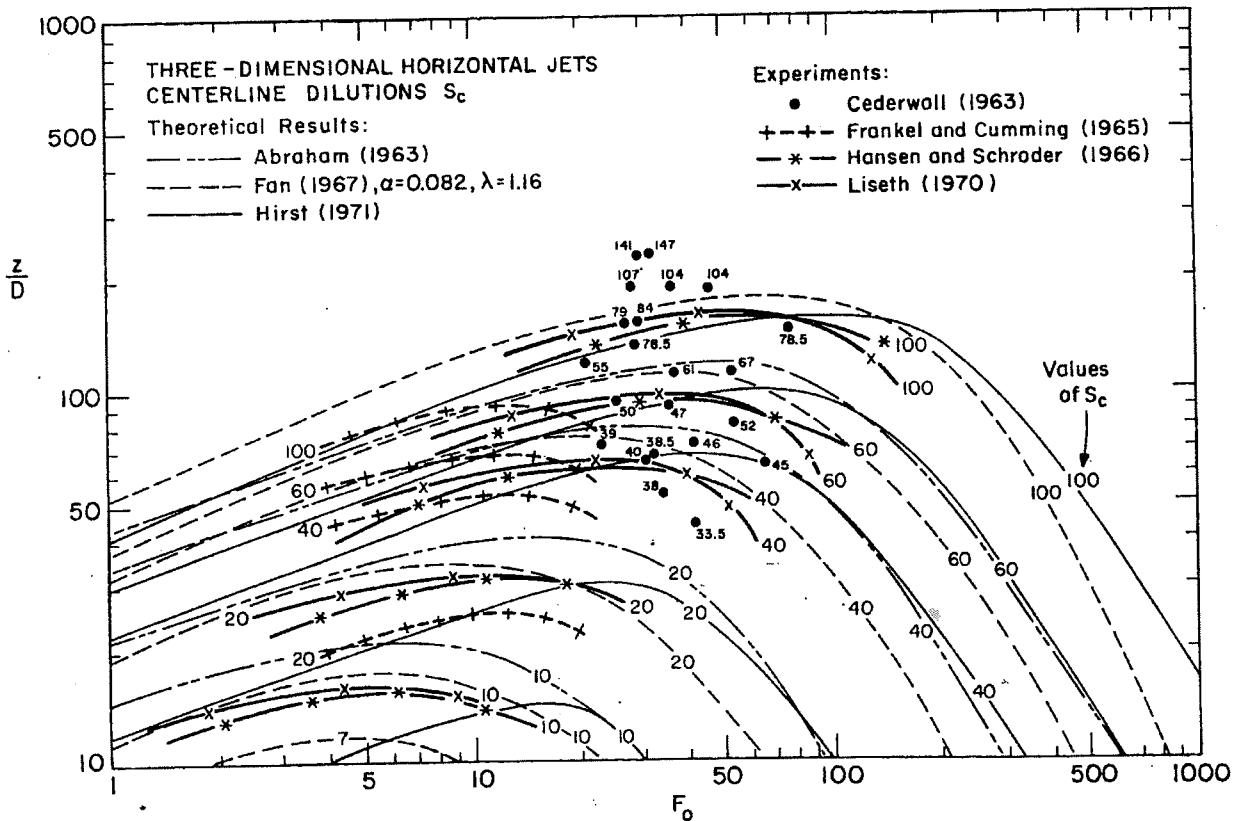


Three-Dimensional Horizontal Jets Centerline Dilutions S_c .
Comparison Between Theoretical Predictions and Experimental Data

STAGNANT HOMOGENEOUS AMBIENT



Three-Dimensional Horizontal Jet Trajectories.
Comparison Between Theoretical Predictions and Experimental Data



Three-Dimensional Horizontal Jets Centerline Dilutions S_c .
Comparison Between Theoretical Predictions and Experimental Data

Dimensional analysis:

Any buoyant jet property e.g. $x = f(z, M_o, J_o, Q_o, D \dots)$

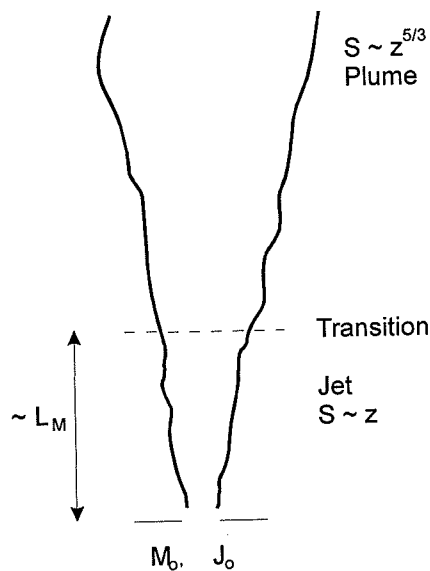
Normalized property $= f\left(\frac{z}{L_M}\right)$

$$L_M = \frac{M_o^{3/4}}{J_o^{1/2}} = \text{momentum length scale}$$

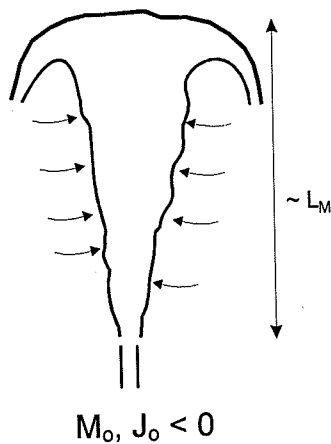
Thus: $DF_o = \left(\frac{4}{\pi}\right)^{1/4} L_M$ correct scaling! $L_M \sim D F_o$

Examples:

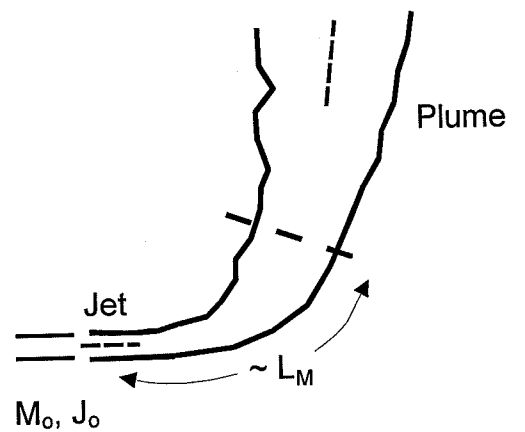
Vertical buoyant jet



Negative buoyant jet

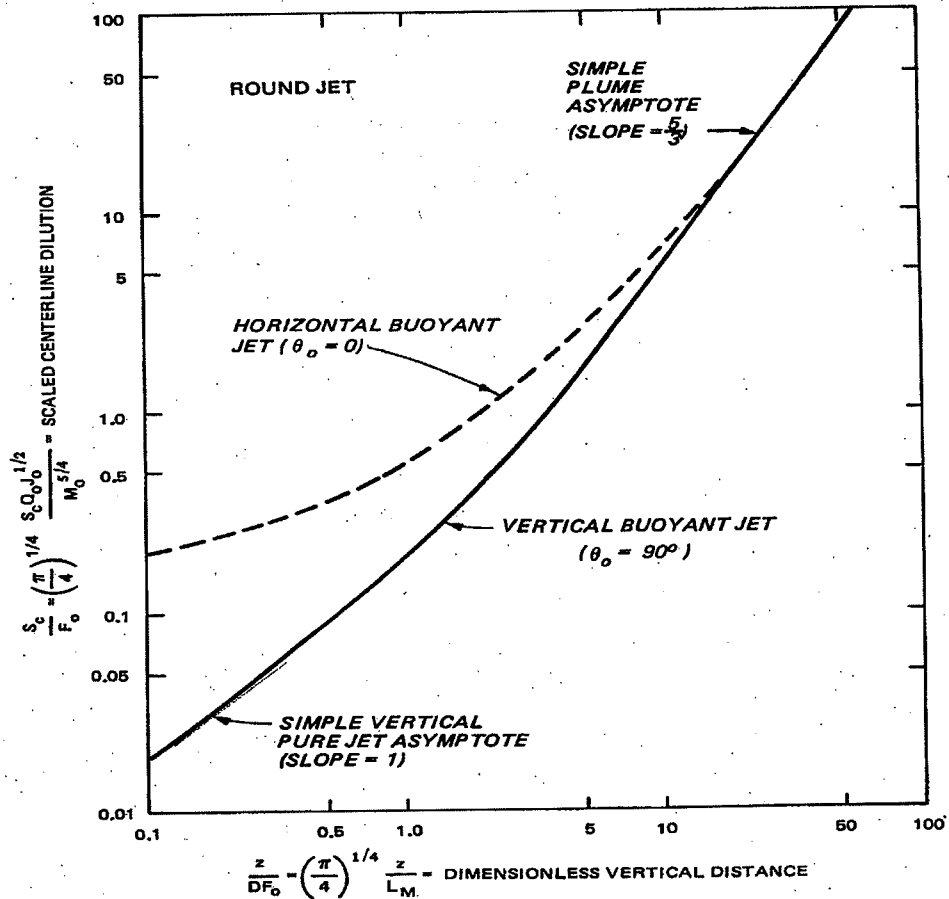
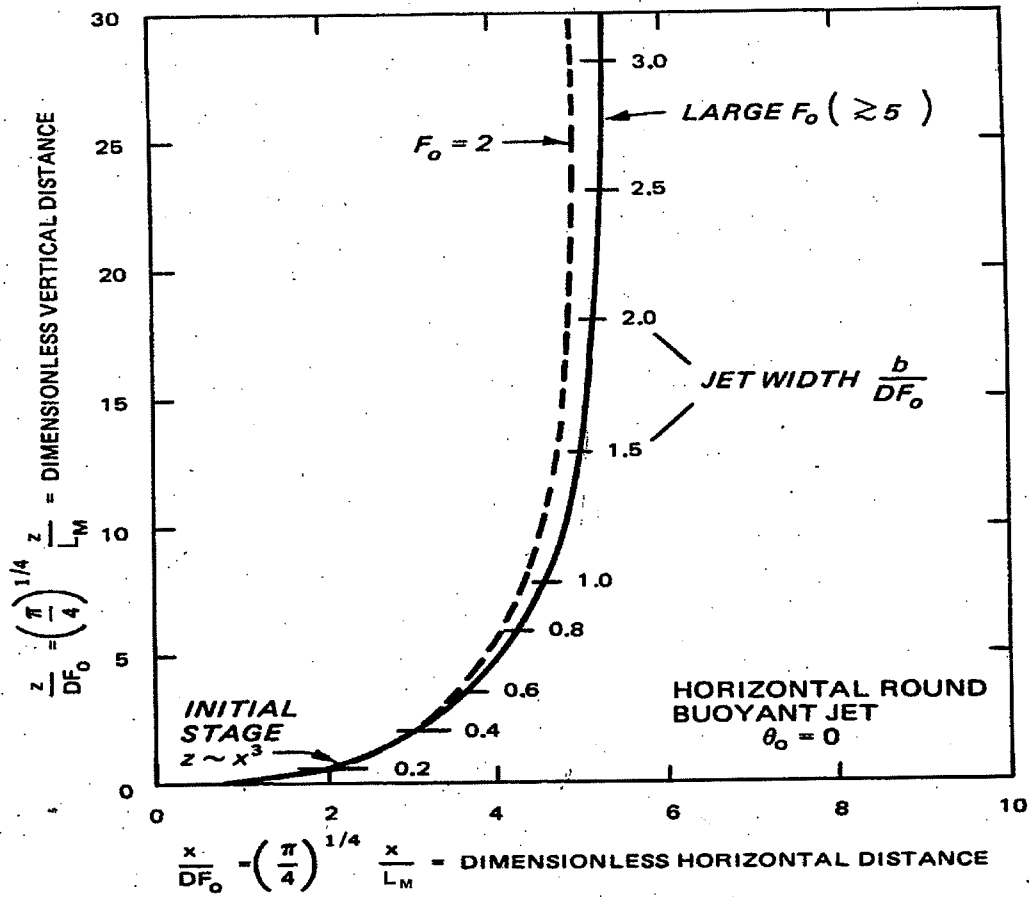


Horizontal buoyant jet



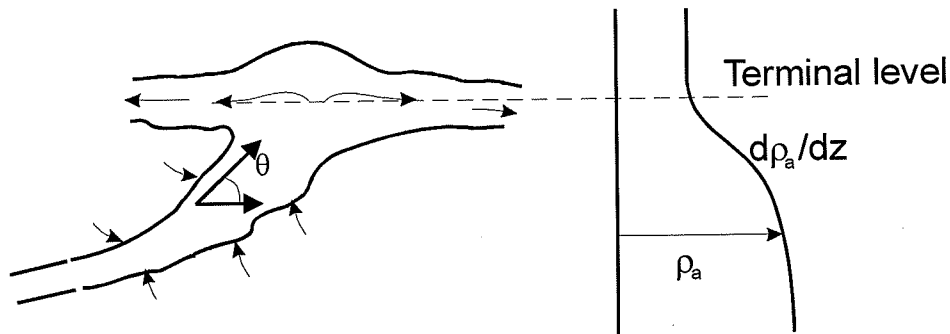
BUOYANT JETS

Stagnant, Unstratified Ambient



Amplifications of the Integral Method

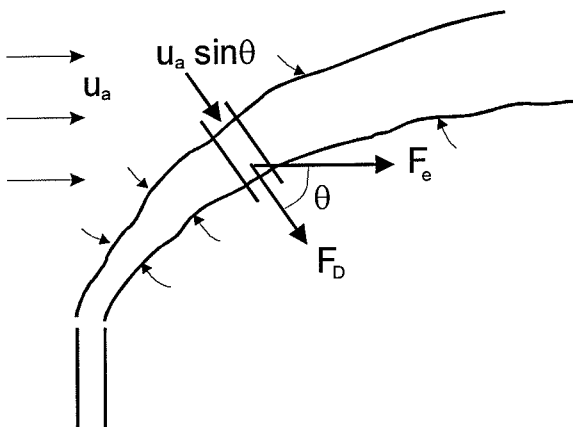
Ambient stratification:



$$(4) \quad \frac{dJ}{ds} = \pi u_c b^2 \frac{d}{dz} \left(\frac{\rho_a g}{\rho_o} \right) \sin \theta = \text{change of buoyancy flux}$$

$$T = \frac{\Delta \rho_o}{D \left(-\frac{d\rho_a}{dz} \right)} \quad \text{Stratification parameter}$$

Crossflow:



Entrainment force =

$$F_e = u_a \frac{dQ}{ds}$$

Drag force =

$$F_D = C_D \frac{(u_a \sin \theta)^2}{2} \quad (2b)$$

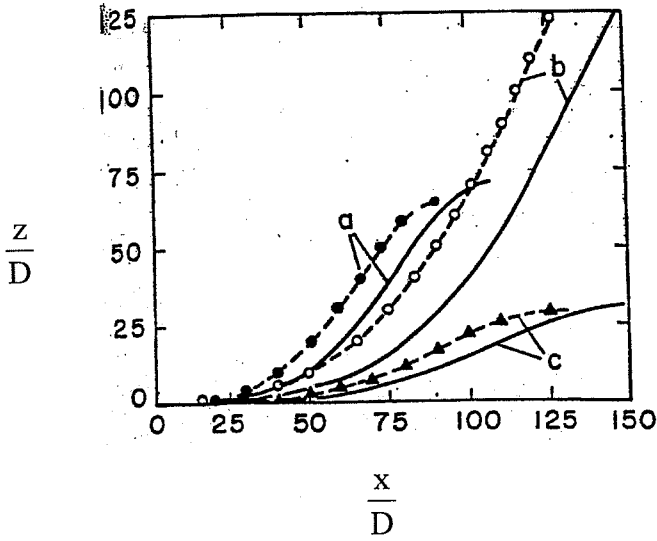
$$(2) \quad \frac{dM}{ds} = \pi \lambda^2 g'_c b \sin \theta + F_e \cos \theta$$

$$(3) \quad \frac{dM \cos \theta}{ds} = F_D \sin \theta + F_e$$

Entrainment \rightarrow additional components: $E = 2\pi b u_c \left(\alpha_1 + \frac{\alpha_2 \sin \theta}{F_l^2} \right) + 2\pi b u_a \alpha_3 \sin \theta \cos \theta$

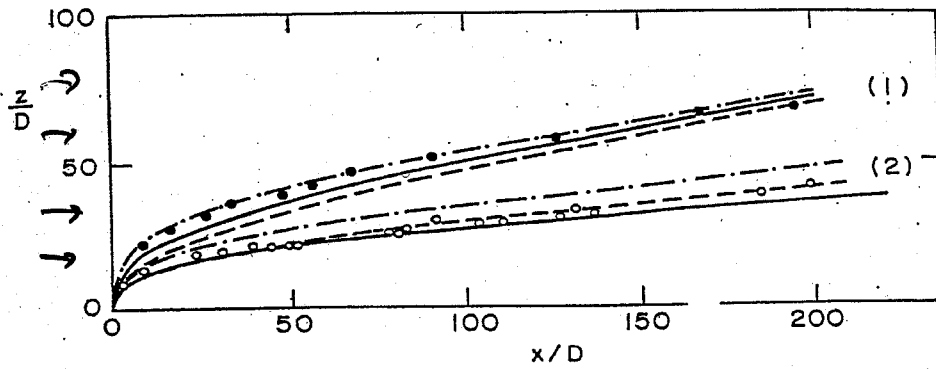
$$R = \frac{u_a}{U_o} \quad \text{Crossflow parameter}$$

Buoyant jet in ambient stratification



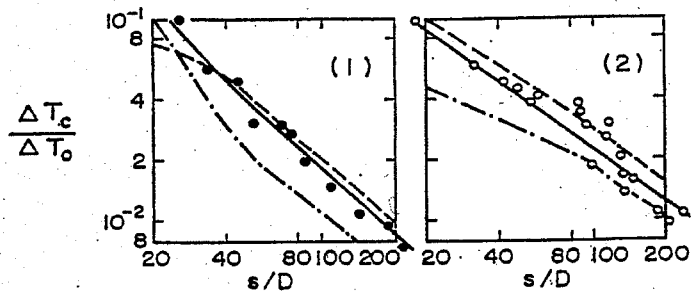
- a: $F_o = 26, T = 1200, \theta_o = 0^\circ$
 - b: $F_o = 40, T = 5000, \theta_o = 0^\circ$
 - ▲ c: $F_o = 60, T = 510, \theta_o = 0^\circ$
 - Fan's model $\alpha = 0.082, \lambda = 1.16$
 - Hirst's model α Eq. (4.4), $\lambda = 1.16$
- Experiments by Fan (1967)

Buoyant jet in crossflow



a) Jet Trajectories

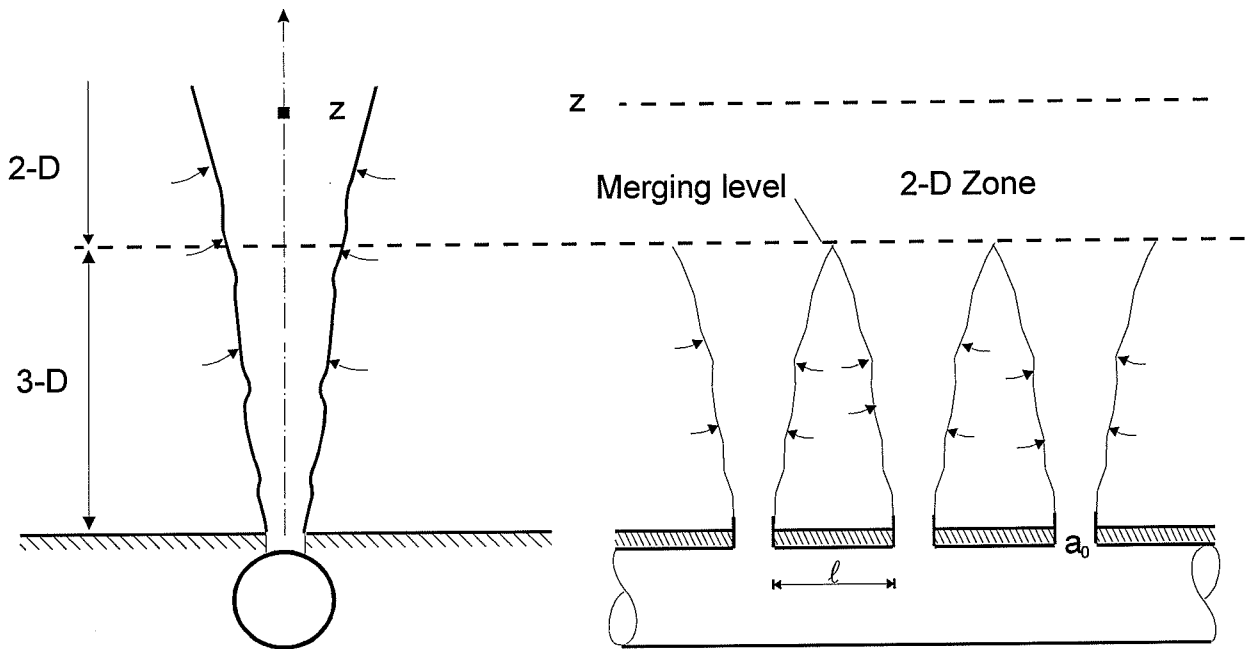
- (1): $F_o = 40, R = 0.063, \theta_o = 90^\circ$
 - (2): $F_o = 40, R = 0.125, \theta_o = 90^\circ$
 - Fan's model, $\alpha = 0.4, C_D = 0.3$ (1), 0.7 (2) (best fit)
 - Abraham's model, α Eq. (4.13), $C_D = 0.4$ (best fit)
 - Hirst's model, no drag
- Experiments by Fan (1967)



b) Centerline temperature (Concentration) Decay
 Buoyant Jets in Crossflow. Comparison of Theoretical Prediction with Experimental Data

Multiple Jets: Multiport diffusor

Analogy to 2-D plume



$$\frac{J_o}{L_D} = j_o = \frac{\Delta\rho_o}{\rho_a} g U_o B \quad \text{buoyancy flux / length} \quad B = \frac{a_o}{l} \quad \text{equivalent slot width}$$

$$q_{c_o} = c_o U_o B$$

Dynamic properties:

$$F_1(b, z, j_o, \dots) = 0$$

$$b = k z$$

$$F_2(u_c, z, j_o, \dots) = 0$$

$$u_c = K j_o^{1/3} = \text{const !}$$

$$F_3(q, z, j_o, \dots) = 0$$

$$q = \bar{K} j_o^{2/3} z$$

Scalar: $q_{c_o} = c_o U_o B \cong c_c u_c b$

$$S = \frac{c_o}{c_c} \sim z \quad \text{Dilution}$$

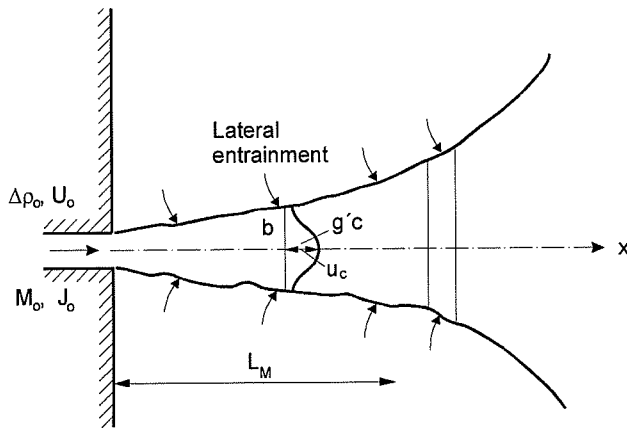
Compare: 3-D plume $S \sim z^{5/3}$

Appropriate length scale for 2-D buoyant jet

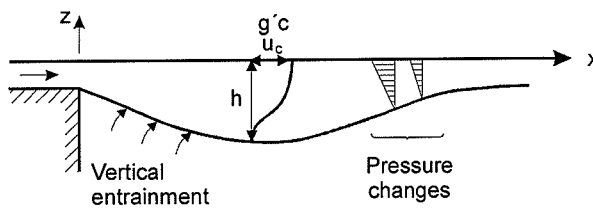
$$l_M = \frac{m_o}{j_o^{2/3}} = \text{jet/plume transition length scale}$$

$$m_o = \frac{M_o}{L_D}, \quad j_o = \frac{J_o}{L_D} \quad L_D = \text{diffusor length}$$

Surface Buoyant Jet



stagnant ambient



$$Q = \int_A u_c f\left(\frac{y}{b}, \frac{z}{h}\right) dA$$

entrainment rate

$$(1) \quad \frac{dQ}{dx} = \underbrace{c_1 \alpha_h u_c h}_{\text{lateral } \alpha_h = \alpha} + \underbrace{c_2 \alpha_v u_c b}_{\text{vertical } \alpha_v = f(F_\ell)} \quad F_\ell = u_c / (g'_c h)^{1/2}$$

e.g. $\alpha = \alpha e^{-5 \cdot F_\ell}$ local buoyancy effect

$$(2) \quad \frac{dM}{dx} = \frac{dP}{dx} \quad P = \int_A \left(\int_{-\infty}^{z/h} \frac{\Delta \rho}{\rho} g d\left(\frac{z}{h}\right) \right) dA \quad \text{buoyant pressure force}$$

$$(3) \quad \frac{db}{dx} = k + \left(\frac{db}{dx}\right)_B \quad u_f = \left(\frac{db}{dt}\right)_B \cong \left(\frac{\Delta \rho_c}{\rho} gh\right)^{1/2} = (g'_c h)^{0.5} \text{ spreading velocity}$$

jet spreading $\approx 0,1$ buoyant spreading $\frac{dx}{dt} \cong u_c$

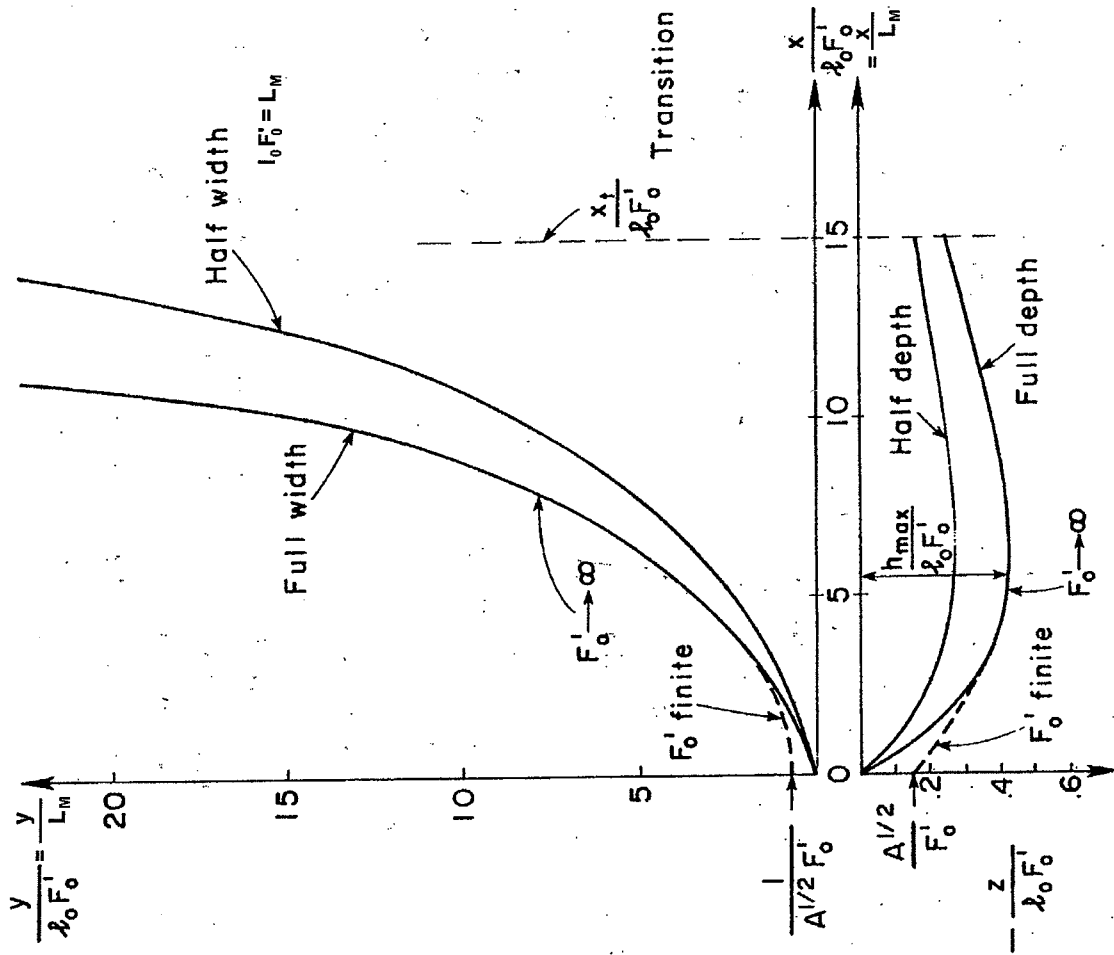
$$\left(\frac{db}{dx}\right)_B = \left(\frac{c_3}{F_\ell}\right) \text{ global buoyancy effect}$$

$$(4) \quad \frac{dJ}{dx} = 0 \quad (\text{no buoyancy loss; e.g. heat loss})$$

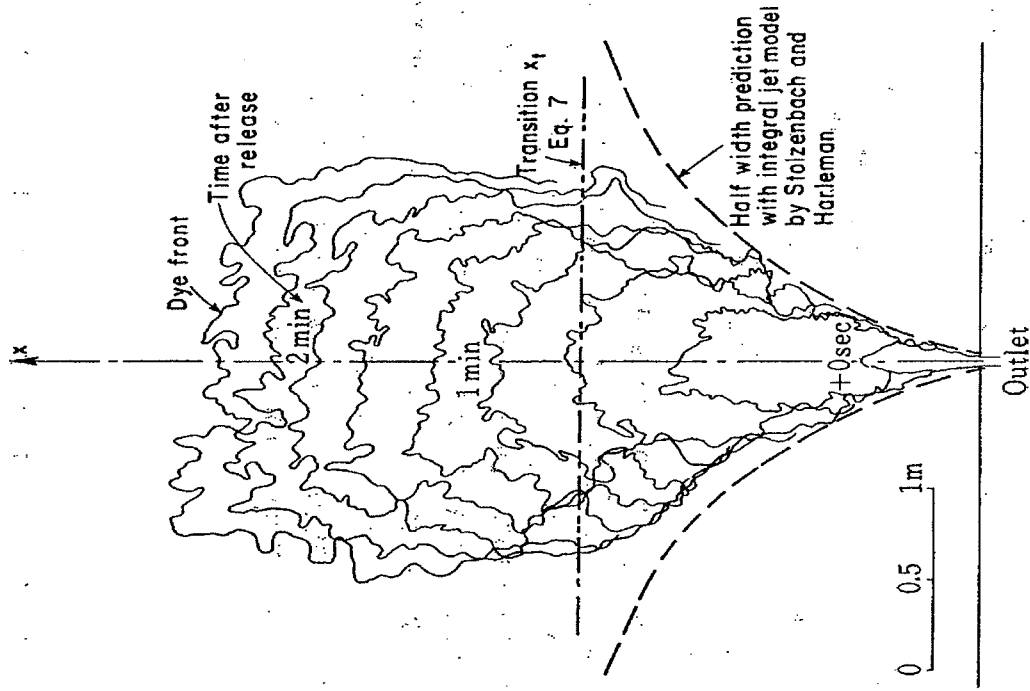
Parameter: F_o

Length Scale: $L_M = \frac{M_o^{3/4}}{J_o^{1/2}}$

Buoyant Surface Jets in Stagnant Ambient

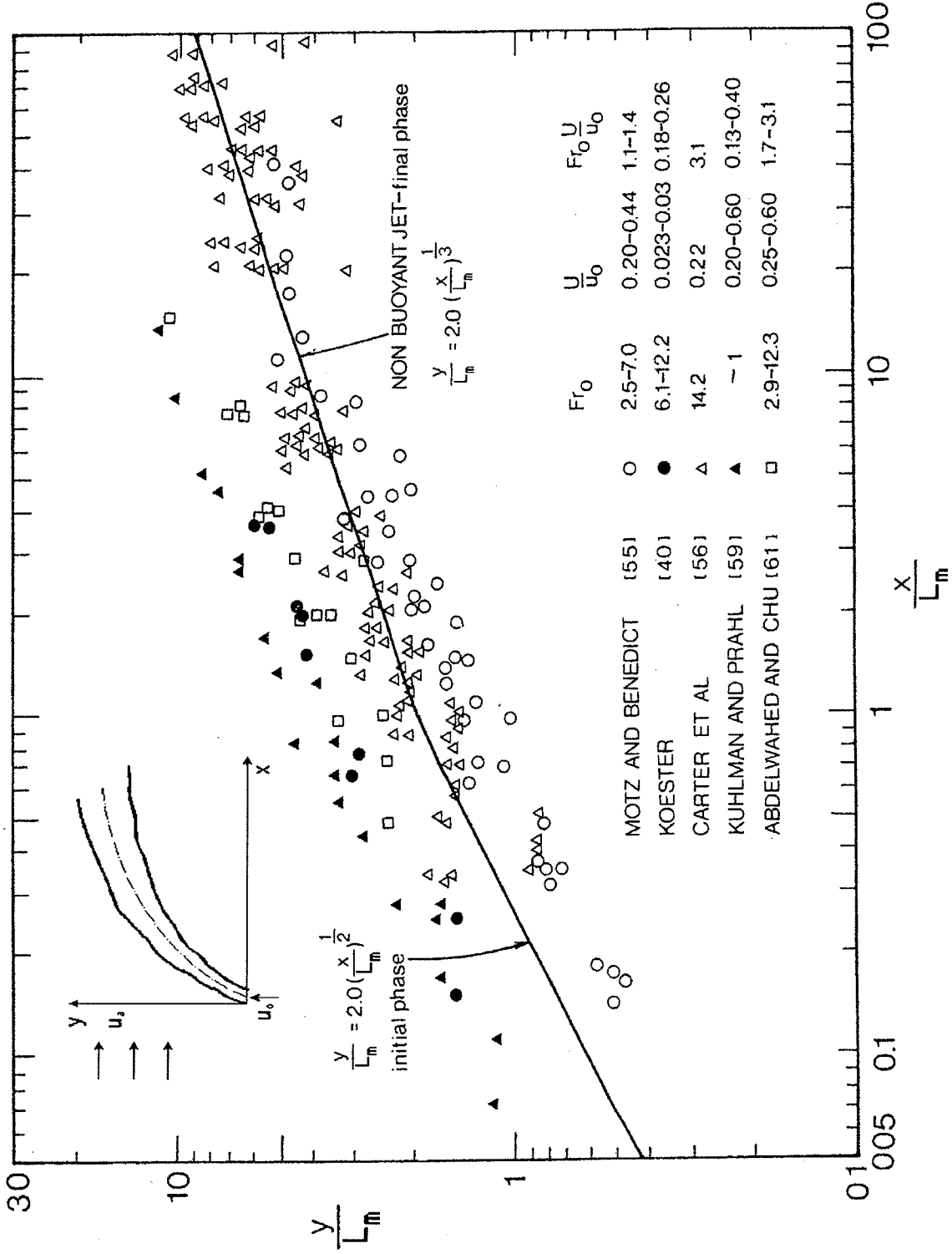


Self-similar Geometry of Buoyant Surface Jets in Deep Receiving Water Computed by Mathematical Model



Unsteady Dye Fronts within Steady Buoyant Surface Jet Showing Nonlinear Spread and Transition. Experiments by Hayashi and Shuto and Mathematical Model Predictions for Jet Half Width

Buoyant Surface Jets in Crossflow



Trajectories of free deflected buoyant jets in a crossflow in comparison to trajectory laws for non-buoyant jets

References: Shallow flow instabilities

- Cantwell, B.J., 1981, "Organized motion in turbulent flow", Ann. Rev. Fluid Mech., 13, 457-515
- Chen, D. and Jirka, G.H., 1995, "Experimental Study of Plane Turbulent Wakes in a Shallow Water Layer", Fluid Dynamics Research, 16, 11-41
- Chen, D. and Jirka, G.H., 1997, "Absolute and Convective Instabilities of Plane Turbulent Wakes in a Shallow Water Layer", J. Fluid Mechanics, 338, 157-172
- Chu, V.H., Wu, J.H. and Khayat, R.E., 1983, "Stability of turbulent shear flows in shallow channel", Proc. XX Congress IAHR, Moscow 3, 128, 133.
- Dracos, T., Giger, M. and Jirka, G.H., 1992, "Plane Turbulent Jets in a Bounded Fluid Layer", J. Fluid Mechanics, 214, 587-614
- Drazin, P.G. and Reid, W.H., 1981, "Hydrodynamic stability", Cambridge University Press, Cambridge
- Giger, M., Dracos, T. and Jirka, G.H., 1991, "Entrainment and Mixing in Plane Turbulent Jets in Shallow Water", J. Hydraulic Research, 29, No.4, 615-643
- Huerre, P. and Monkewitz, P.A., 1990, "Local and global instabilities in spatially developing flows", Ann. Rev. Fluid Mech. 22, 473-537
- Huerre, P. and Monkewitz, P.A., 1990, "Local and global instabilities in spatially developing flows", Ann. Rev. Fluid Mech., 22, 472-537
- Jirka, G.H., 1994, "Shallow Jets", in Recent Advances in the Fluid Mechanics of Turbulent Jets and Plumes, P.A. Davies and M.J. Valente Neves (Ed.s), Kluwer Academic Publishers, Dordrecht
- Monkewitz, P.A., 1988, "The absolute and convective nature of instability in two-dimensional wakes at low Reynolds number", Phys Fluid 31 (5), 999-1006
- Oertel, H., 1990, "Wakes behind blunt bodies", Ann. Rev. Fluid Mech. 22, 530-564
- Van Dyke, M., 1982, "An Album of Fluid Motion", Parabolic Press, Stanford, CA
- Wolanski, E., Imberger, J. and Heron, M.L., 1984, "Island wakes in shallow and coastal waters", J. Geophys. Res. 89 (C6), 10553-10569