

# Methods for Obtaining Russell-Saunders Term Symbols from Electronic Configurations

A thorough understanding of the concepts of energy terms, levels, and microstates is paramount for students engaging in the study of quantum mechanics, spectroscopy, and the crystal or ligand field theories of bonding in transition metal complexes. In teaching courses of physical or inorganic chemistry to our undergraduate students we have found that few texts present a scheme that may be used to derive the appropriate energy level diagram for configurations more complex than the  $p^2$  case. Most of the available texts recommend forming all possible arrangements of electrons in the available orbitals and discarding those which violate either the Pauli principle or are identical with one another by the indistinguishability of electrons. The problem of equivalent or nonequivalent electrons only adds to the confusion. The case illustrated is most often  $p^2$  but we believe that this is one of the few multielectron cases that can be handled by the above technique without extreme confusion and tedium. The purpose of this article is to illustrate and comment on several other procedures that are available and may be used for the more complex electron configurations.

Of the techniques described in current texts the method advocated by Douglas and McDaniel (1) leads to a maximum of understanding and a minimum of frustration. Surprisingly, this method has not appeared in any of the more recent inorganic (2-4) or physical texts. When this method has been thoroughly explained, our students have been capable of generating the correct energy levels for electron configurations as complex as  $d^4$ ,  $d^5$ ,  $f^2$ ,  $f^3$  and mixed configurations such as  $d^2s^1$  and  $d^1f^2$ . The application of this method is described under the heading Method I. It should be noted that while the original ideas for this method were presented by Douglas and McDaniel (1), the treatment offered here is fuller and in more detail.

Another procedure illustrated by E. R. Tuttle (5) for obtaining terms for configurations involving equivalent electrons is about as convenient as that of Douglas and McDaniel's but its weakness lies in the fact that a set of rules (the origin of these rules is not immediately obvious to the beginning student) must be adopted and a process of repetitively applying these rules leads to the proper terms for the configuration of interest. A detailed description of this treatment may be obtained by writing this author.

A third technique involving generating functions derived via group theory has been developed by Curl and Kilpatrick (6). The main disadvantage lies in an excess of algebraic manipulations. Without a knowledge of group theory the origin of the generating functions and their relation to the problem at hand cannot be demonstrated. To obtain the proper term symbols from a given configuration, group theory is not required. A fuller treatment of this procedure may also be obtained from this author.

As Phillips and Williams (7) point out, additional confusion exists over what is meant by the use of the words term, state, level, and microstate. We advocate the use of the terminology given in Table 1, when the Russell-Saunders coupling scheme is applicable. In order to alleviate a confusing situation, it is strongly recommended that this terminology be adopted by all who use these descriptions.

As a means of illustrating the application of the pre-

**Table 1. Terminology Used in the Generation of Energy Levels for Multielectron Atoms**

Nomenclature	Quantum numbers specified <sup>a</sup>	Symbol
Configuration	$n$ and $l$ (the number of electrons $N$ is also specified)	$n l^N$
Term	$n, L, S$	$2S + 1 L$
Level	$n, L, S, J$	$2S + 1 L_J$
Microstate	$n, L, S, J, M_J$	$2S + 1 L_J (M_J)$

<sup>a</sup>The usual convention of lower case letters representing single electron quantum numbers and upper case letters representing those for a multielectron atom has been used.

**Table 2. Possible  $M_S$  Values for Groupings of Arrangements with Zero to Seven Unpaired Electrons<sup>a</sup>**

Number of unpaired electrons in group ( $\beta$ )	$M_S^b$ values for group	$2\beta^c$
0	0	1
1	$+1/2, -1/2$	2
2	$1, 0(2), -1$	4
3	$+3/2, +1/2(3), -1/2(3), -3/2$	8
4	$2, 1(4), 0(6), -1(4), -2$	16
5	$+5/2, +3/2(5), +1/2(10), -1/2(10), -3/2(5), -5/2$	32
6	$+3, 2(6), 1(15), 0(20), -1(15), -2(6), -3$	64
7	$+7/2, +5/2(7), 3/2(21), 1/2(35), -1/2(35), -3/2(21), -5/2(7), -1/2$	128

<sup>a</sup>The table can be generated by adding  $\pm 1/2$  to the possible  $M_S$  values for an arrangement with  $\beta$  unpaired electrons to give the possible  $M_S$  values for grouping with  $\beta + 1$  unpaired electrons.

<sup>b</sup>The number in parentheses indicates the number of times that particular  $M_S$  value will appear for a configuration having the specified  $\beta$  unpaired electrons.

<sup>c</sup>The quantity  $2\beta$  represents the total number of arrangements or  $M_S$  values that can be written for  $\beta$  unpaired electrons.

**Table 3. Grouping of Arrangements with  $S = 1/2$**

$M_L$						$M_L$							
$m_l =$	2	1	0	-1	-2	$\Sigma m_l$	$m_l =$	2	1	0	-1	-2	$\Sigma m_l$
X	/					5				X	/		-1
X		/				4				X		/	-2
X			/			3		/		X			0
X				/		2		/		X			-1
/	X					4			/	X			-2
	X	/				2				X	/		-4
	X		/			1		/		X			-2
	X			/		0		/		X			-3
/		X				2			/		X		-4
	/	X				1				/	X		-5

viously described method we chose the  $nd^3$  case as an example. For three  $d$  electrons the total degeneracy of microstates is given by

$$T.D. = \frac{(4l + 2)!}{(4l + 2 - N)!N!}$$

where  $N$  is the number of electrons in the configuration and  $l$  is the secondary quantum number. For  $nd^3$  we have  $l = 2$  for a  $d$  electron,  $N = 3$ , and  $T.D. = 10!/(7! \times 3!) = 120$ . There are 120 ways of arranging three  $d$  electrons in the five degenerate  $d$  orbitals. The problem is to sort out the energy terms contained within these 120 arrangements.

## Method I

For three  $d$  electrons there are two possible spin states, characterized by  $S = 1/2$  and  $S = 3/2$ . Electron pairs are



### The Major Array that Contains the 120 Arrangements

The array is constructed by realizing that each  $M_L$  value that occurs in arrangements with but one unpaired electron has two possible  $M_S$  values while each  $M_L$  value that occurs in arrangements with three unpaired electrons has eight  $M_S$  values.

$M_L$					
5		1	1		
4		2	2		
3	1	4	4	1	
2	1	6	6	1	
1	2	8	8	2	
0	2	8	8	2	
-1	2	8	8	2	
-2	1	6	6	1	
-3	1	4	4	1	
-4		2	2		
-5		1	1		
	$3/2$	$1/2$	$-1/2$	$-3/2$	$M_S$

Within this major array are 120 allowed arrangements of three electrons in the five  $d$  orbitals. Alternatively, we can view the major array as containing several subarrays where each subarray defines a term. Since  $L$  and  $S$  are specified for a term, each term is the equivalent of an array of  $M_L$  and  $M_S$  values that arise from the  $L$  and  $S$  values of the term. For instance, the term  $^3P$  ( $L = 1$ ,  $S = 1$ ) is equivalent to the array

$M_L$				
1	1	1	1	
0	1	1	1	
-1	1	1	1	
	1	0	-1	$M_S$

The  $^3P$  term is three-fold ( $2S + 1$ ) spin degenerate and three-fold ( $2L + 1$ ) orbitally degenerate and hence has a total degeneracy of nine. Nine microstates or nine arrangements of electrons exist.

The problem that exists after the construction of the major array is simply one of determining what subarrays (or terms) are contained within the major array. The simplest procedure is one of eliminating all  $M_L$  and  $M_S$  values (a subarray) that belong to a particular term, thus producing a new major array. This process is repeated until the major array vanishes.

One can always pick the term of maximum  $L$  from the major array by realizing that  $L$  for this term is given by the maximum value of  $M_L$  in the major array. This must be true since the  $M_L$  values of this term are related to  $L$  by  $M_L = L, L - 1, \dots, 0, \dots, -L$  ( $M_L$  takes all values from

$+L$  through zero to  $-L$ ). Once  $L$  or  $M_L(\max)$  is established for the term the  $S$  value is given by the maximum value for  $M_S$  with this same  $M_L(\max)$ . This also follows since  $M_S = S, S - 1, \dots, 0, \dots, -S$ ; hence  $S$  for the term under consideration is  $M_S(\max)$ .

To illustrate by use of the major array generated for the  $nd^3$  configuration we note that  $M_L(\max) = 5$ ; hence  $L = 5$  or an  $H$  type term is present. The maximum value in the array for  $M_S$  is  $+1/2$  if  $M_L$  is still at its maximum value of five; therefore  $M_S(\max) = S = +1/2$ . The conclusion is that the term  $^2H$  is present in the major array.

The subarray of  $M_L$  and  $M_S$  values generated by the  $^2H$  term is subtracted from the major array. The new major array generated contains the  $^2G$  term, which must be removed from the new array. The procedure is repeated until the major array vanishes. The figure illustrates this procedure diagrammatically. Results indicate that the  $nd^3$  configuration contains the terms  $^2H$ ,  $^2G$ ,  $^4F$ ,  $^2F$ ,  $^2D(\text{twice})$ ,  $^4P$ , and  $^2P$ .

### Conclusion

Herein have been presented three alternative methods for deriving the Russell-Saunders terms for any electronic configurations. The advantages and disadvantages of each have been mentioned at the beginning. When only the ground term for a configuration is desired, it can be obtained by a straightforward and simple procedure. A recent article in *this Journal* (8) illustrates the process and it will not be described here.

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